

# ODEs

Prateek Sharma ([prateek@physics.iisc.ernet.in](mailto:prateek@physics.iisc.ernet.in))

Office: D2-08

# Reduction to 1<sup>st</sup> order ODEs

$$\frac{d^2 y}{dx^2} + q(x) \frac{dy}{dx} = r(x)$$

$$\frac{dy}{dx} = z(x)$$

$$\frac{dz}{dx} = r(x) - q(x)z(x)$$

new variables introduced s.t.  
equations are well-behaved

generic set of  $N$  coupled first-order ODEs:

$$\frac{dy_i(x)}{dx} = f_i(x, y_1, \dots, y_N), \quad i = 1, \dots, N$$

# BVPs & IVPs

Boundary conditions are algebraic conditions on the values of the functions  $y_i$  at the boundaries. They can be satisfied at discrete specified points. Two broad categories:

*initial value problems* all the  $y_i$  are given at some starting value  $x_s$ , and it is desired to find the  $y_i$ 's at some final point  $x_f$ , or at some discrete list of points. e.g., planetary orbits with given initial conditions

*two-point boundary value problems* boundary conditions are specified at more than one  $x$ ; some of the conditions will be specified at  $x_s$  and the remainder at  $x_f$ ; e.g., stellar structure: core density, temperature at inner boundary and radius, luminosity at outer boundary. in general more difficult.

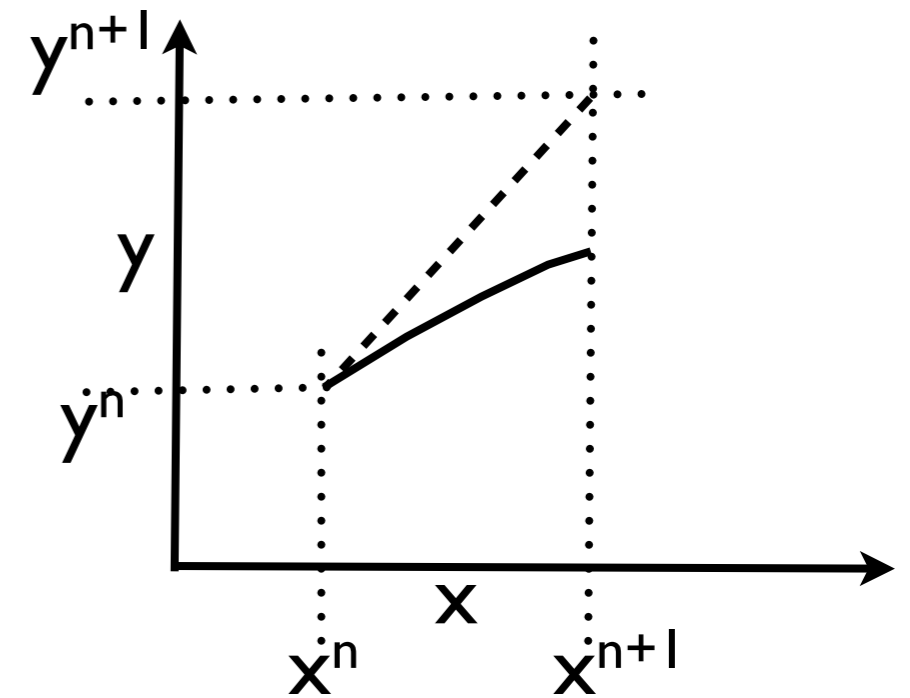
lets consider IVPs for now

# Forward Euler

$$\frac{y_i^{n+1} - y_i^n}{\Delta x} = f_i(x^n, y_1^n, \dots, y_N^n); i = 1, \dots, N$$

$$y^{n+1} = y^n + \frac{dy}{dx} \Delta x + \frac{d^2 y}{dx^2} \frac{\Delta x^2}{2} + \dots$$

$$\frac{dy}{dx} + f' \frac{\Delta x}{2} = f \quad \text{modified eq. 1}^{\text{st}} \text{ order accurate}$$



numerically unstable for large  $\Delta x$ ! e.g.,  $dy/dx = -ay$ ,  $y(0) = 1 \Rightarrow y = e^{-ax}$

$$\text{FE: } y^{n+1} = (1 - ah)^{n+1} y^0$$

$|1 - ah| < 1$  for stability!

# Backward Euler

$$\frac{y_i^{n+1} - y_i^n}{\Delta x} = f_i(x^{n+1}, y_1^{n+1}, \dots, y_N^{n+1}); i = 1, \dots, N$$

also 1<sup>st</sup> order accurate; this is an implicit eq.  
to be solved, e.g., via Newton-Raphson  
unconditionally stable!

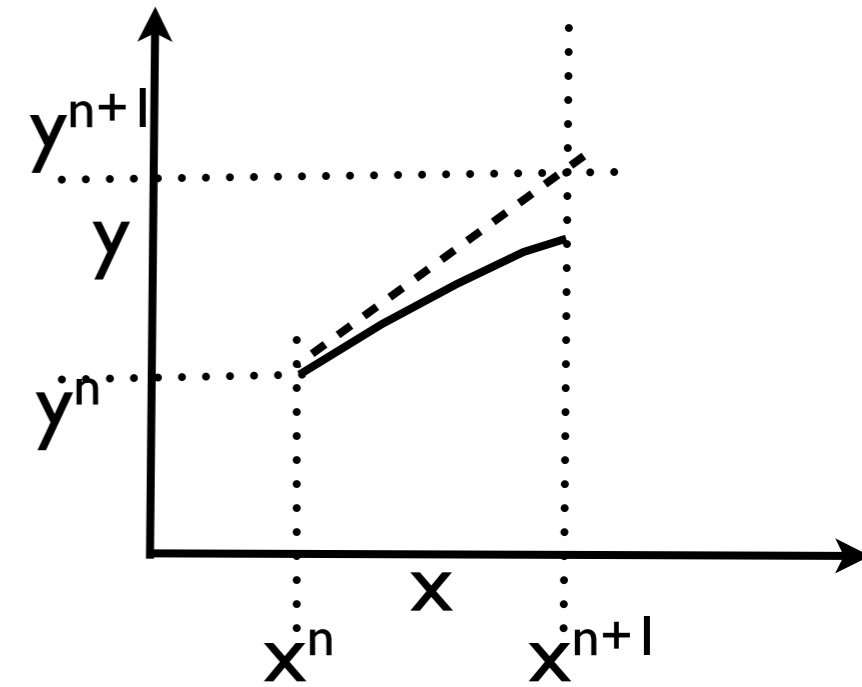


Figure 2: Numerical instability in the FE method for  $dt > 0.2$ ; IVP is  $dy/dx = -10y$ ,  $y(0)$

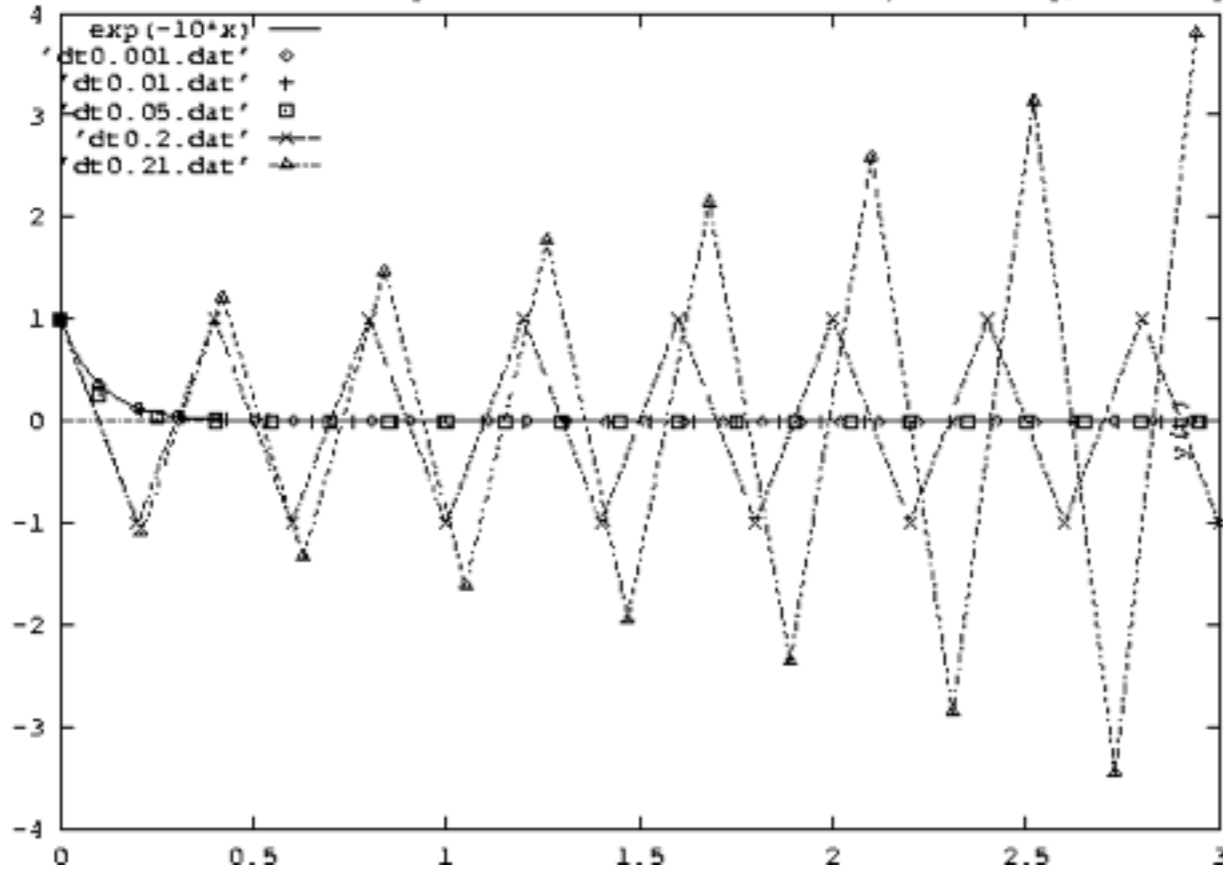
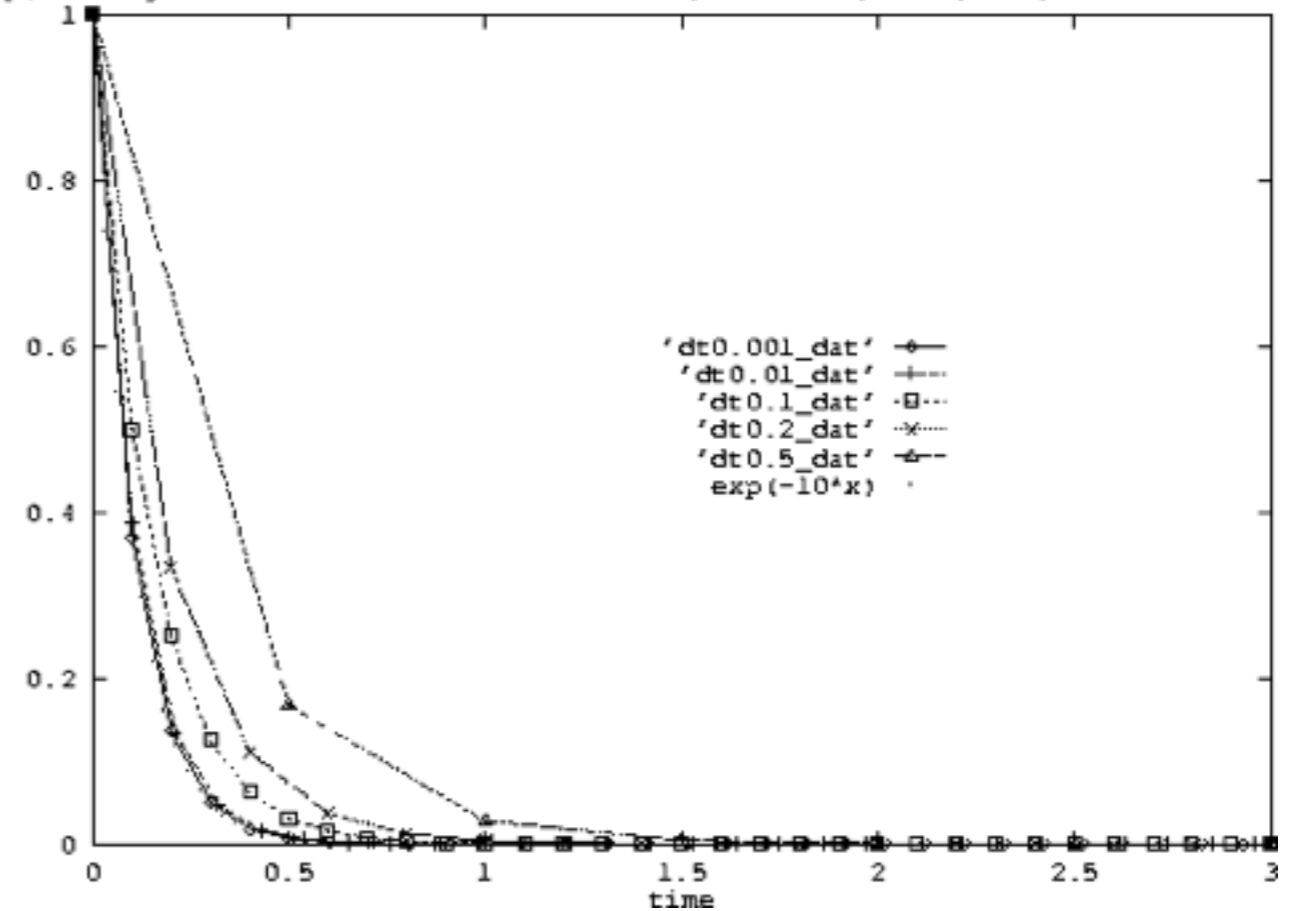


Figure 4: The backward Euler method,  $dt=0.001, 0.01, 0.1, 0.2$  and  $0.5$



# Runge-Kutta Schemes

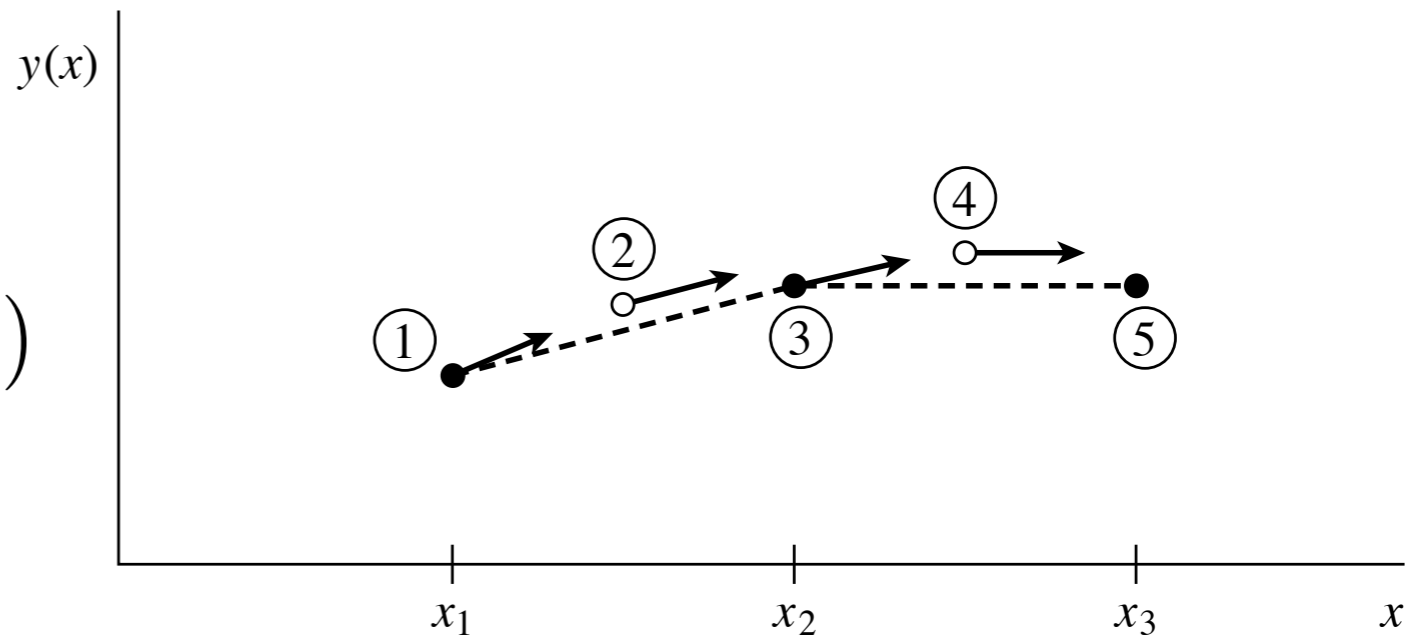
Euler is only 1<sup>st</sup> order accurate!

center the derivative at  $n+1/2$ : mid-point/RK2

$$k_1 = hf(x_n, y_n)$$

$$k_2 = hf\left(x_n + \frac{1}{2}h, y_n + \frac{1}{2}k_1\right)$$

$$y_{n+1} = y_n + k_2 + O(h^3)$$



2<sup>nd</sup> order accurate!

Can be easily shown via Taylor series

# RK4

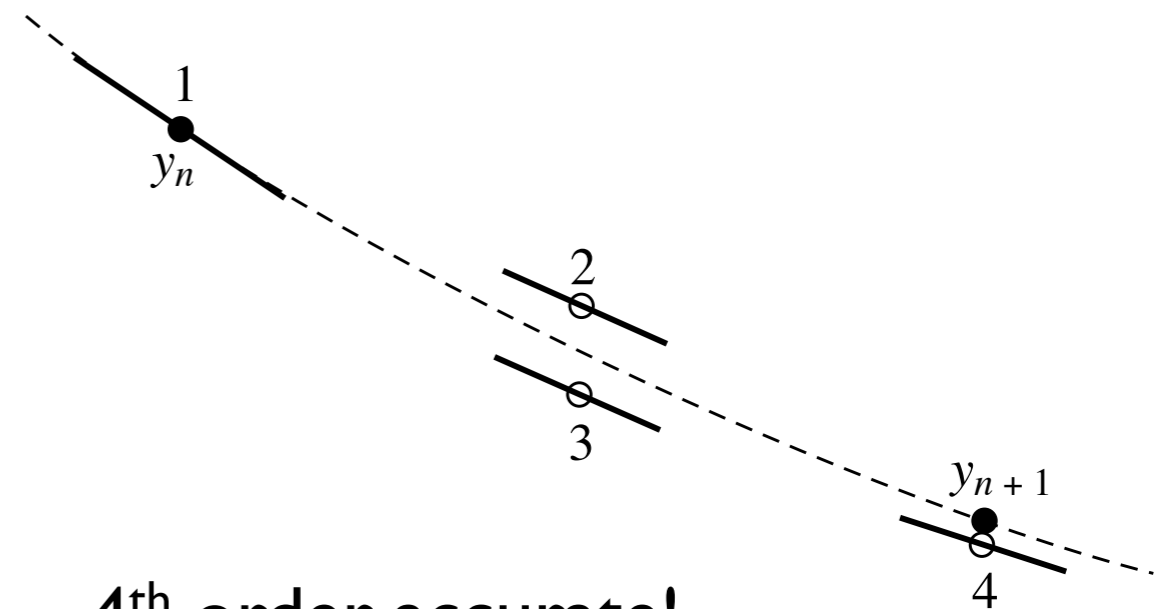
$$k_1 = hf(x_n, y_n)$$

$$k_2 = hf\left(x_n + \frac{h}{2}, y_n + \frac{k_1}{2}\right)$$

$$k_3 = hf\left(x_n + \frac{h}{2}, y_n + \frac{k_2}{2}\right)$$

$$k_4 = hf(x_n + h, y_n + k_3)$$

$$y_{n+1} = y_n + \frac{k_1}{6} + \frac{k_2}{3} + \frac{k_3}{3} + \frac{k_4}{6} + O(h^5)$$



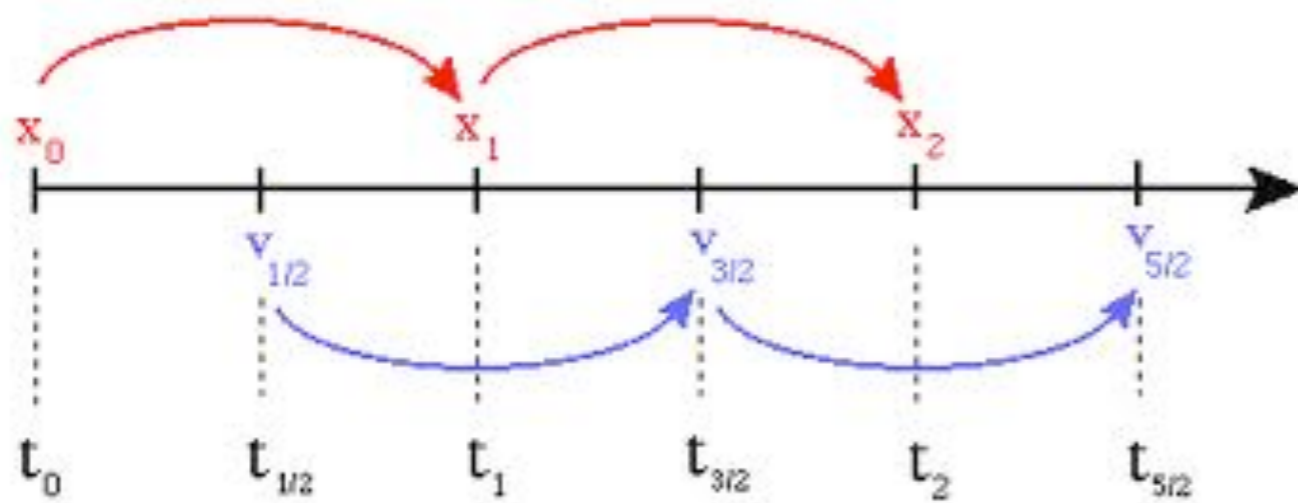
4<sup>th</sup> order accurate!

most common ODE solver; while more accurate, also requires twice the number of function evaluations

# Leap-frog Method

$\ddot{\mathbf{x}} = F(\mathbf{x})$  Newton's 2<sup>nd</sup> law: e.g., planetary orbits, molecular dynamics, etc.

$\dot{\mathbf{v}} = F(\mathbf{x}), \dot{\mathbf{x}} \equiv \mathbf{v}$       $\ddot{\mathbf{x}} = -\nabla V(\mathbf{x})$       $E(\mathbf{x}, \mathbf{v}) = \frac{1}{2}|\mathbf{v}|^2 + V(\mathbf{x})$   
 time-reversible, conservative,  
 Hamiltonian system



$$\mathbf{x}_i = \mathbf{x}_{i-1} + \mathbf{v}_{i-1/2} \Delta t,$$

$$\mathbf{a}_i = F(\mathbf{x}_i)$$

$$\mathbf{v}_{i+1/2} = \mathbf{v}_{i-1/2} + \mathbf{a}_i \Delta t,$$

$$\mathbf{x}_{i+1} = \mathbf{x}_i + \mathbf{v}_i \Delta t + \frac{1}{2} \mathbf{a}_i \Delta t^2,$$

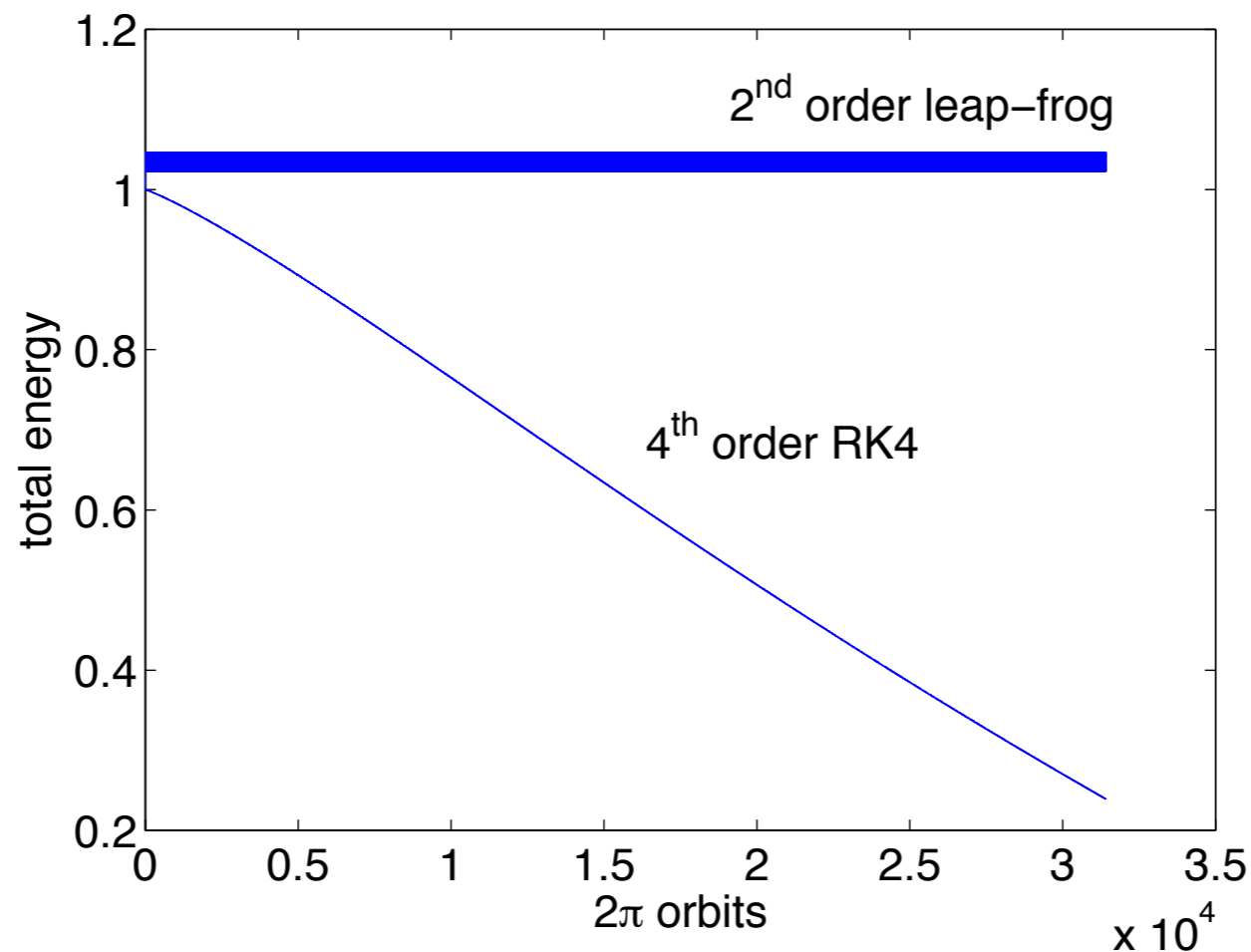
$$\mathbf{v}_{i+1} = \mathbf{v}_i + \frac{1}{2} (\mathbf{a}_i + \mathbf{a}_{i+1}) \Delta t.$$



LF/Verlet very useful where we want to conserve energy, prevent secular errors, not just the formal one-step error  
RK, etc. lead to non-conservation of energy; energy drifts!

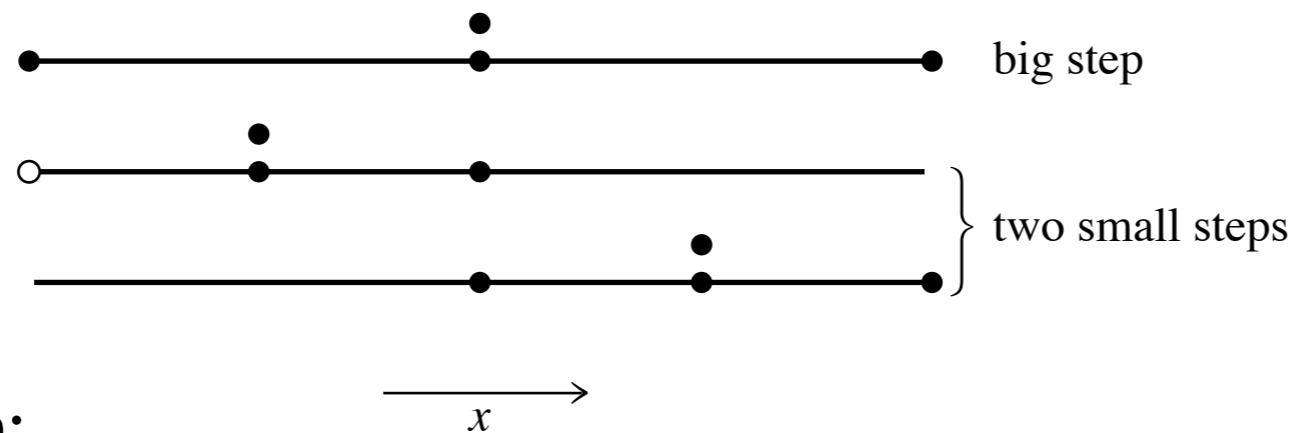
Cons: dt must be the same for time-reversibility, etc. lower order; RK, etc. fine for short-term evolution

RK, LF are explicit schemes and thus require  $dt < 1/\omega_0$  for stability



# Adaptive step-size

stepsize ( $h$ ) is chosen to achieve some pre-specified accuracy; algorithm should give an estimate of truncation error



Taylor expansion:

$$y(x + 2h) = y_1 + (2h)^5 \phi + O(h^6) + \dots$$

$$y(x + 2h) = y_2 + 2(h^5) \phi + O(h^6) + \dots$$

$\Delta \equiv y_2 - y_1$  a good measure of truncation error

$$h_0 = h_1 \left| \frac{\Delta_0}{\Delta_1} \right|^{0.2} \leftarrow \text{desired accuracy} \quad \text{must use an adaptive stepsize}$$

# Embedded RK

Cash-Karp Parameters for Embedded Runge-Kutta Method								
$i$	$a_i$	$b_{ij}$					$c_i$	$c_i^*$
1							$\frac{37}{378}$	$\frac{2825}{27648}$
2	$\frac{1}{5}$	$\frac{1}{5}$					0	0
3	$\frac{3}{10}$	$\frac{3}{40}$	$\frac{9}{40}$				$\frac{250}{621}$	$\frac{18575}{48384}$
4	$\frac{3}{5}$	$\frac{3}{10}$	$-\frac{9}{10}$	$\frac{6}{5}$			$\frac{125}{594}$	$\frac{13525}{55296}$
5	1	$-\frac{11}{54}$	$\frac{5}{2}$	$-\frac{70}{27}$	$\frac{35}{27}$		0	$\frac{277}{14336}$
6	$\frac{7}{8}$	$\frac{1631}{55296}$	$\frac{175}{512}$	$\frac{575}{13824}$	$\frac{44275}{110592}$	$\frac{253}{4096}$	$\frac{512}{1771}$	$\frac{1}{4}$
	$j =$	1	2	3	4	5		

$$k_1 = hf(x_n, y_n)$$

$$k_2 = hf(x_n + a_2h, y_n + b_{21}k_1)$$

...

$$k_6 = hf(x_n + a_6h, y_n + b_{61}k_1 + \dots + b_{65}k_5)$$

$$y_{n+1} = y_n + c_1k_1 + c_2k_2 + c_3k_3 + c_4k_4 + c_5k_5 + c_6k_6 + O(h^6)$$

$$y_{n+1}^* = y_n + c_1^*k_1 + c_2^*k_2 + c_3^*k_3 + c_4^*k_4 + c_5^*k_5 + c_6^*k_6 + O(h^5)$$

$$\Delta \equiv y_{n+1} - y_{n+1}^* = \sum_{i=1}^6 (c_i - c_i^*)k_i$$

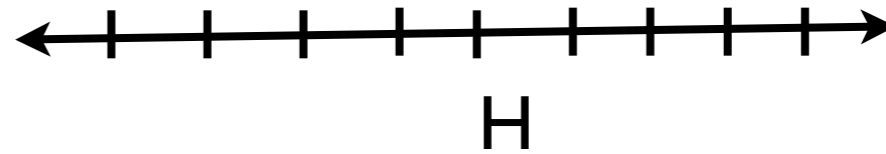
$$h_0 = h_1 \left| \frac{\Delta_0}{\Delta_1} \right|^{0.2}$$

1/2 the no. of fn. evals.!

recall truncation error,  $\Delta_0$  can't be too small

# Modified Midpoint

$$h = H/n$$



$$z_0 \equiv y(x)$$

$$z_1 = z_0 + hf(x, z_0)$$

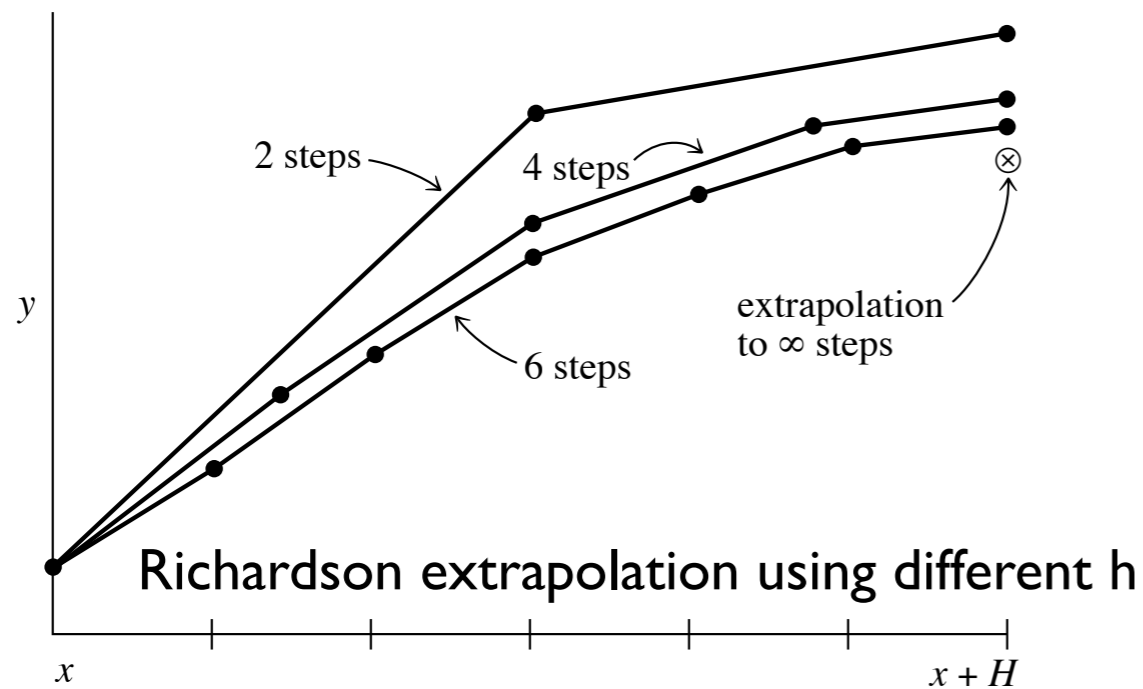
$$z_{m+1} = z_{m-1} + 2hf(x + mh, z_m) \quad \text{for } m = 1, 2, \dots, n-1$$

$$y(x + H) \approx y_n \equiv \frac{1}{2}[z_n + z_{n-1} + hf(x + H, z_n)] \quad \text{second order, with 1 derivative evaluation per h}$$

$$y_n - y(x + H) = \sum_{i=1}^{\infty} \alpha_i h^{2i} \quad y(x + H) \approx \frac{4y_n - y_{n/2}}{3} \quad \text{4th order accurate, only 1.5 fn. evaluations!}$$

apply only for ODEs containing smooth functions

# Bulirsch-Stoer



apply only for ODEs containing smooth functions  
for very high accuracy

RK4 w. adaptive stepsize for non-smooth fns.

rational vs. polynomial extrapolation of error

$$y_n - y(x + H) = \sum_{i=1}^{\infty} \alpha_i h^{2i}$$

use a method with even terms in error

$$n = 2, 4, 6, 8, 10, 12, 14, \dots, [n_j = 2j], \dots$$

sub-intervals; for each n obtain  
approximation & error estimate  
go to higher n if error large

$$H_k = H \left( \frac{\epsilon}{\epsilon_{k+1,k}} \right)^{1/(2k+1)}$$

big stepsize should be small enough, s.t.,  
kth column error is smaller than specified

# Stiff Equations

$$u' = 998u + 1998v$$

$$v' = -999u - 1999v$$

$$u = 2e^{-x} - e^{-1000x}$$

$$v = -e^{-x} + e^{-1000x}$$

vanishingly small, yet determines h!

if there are widely separated timescales in the problem

*implicit* is the way to go

$$u(0) = 1 \quad v(0) = 0$$

$$\mathbf{y}' = \mathbf{f}(\mathbf{y})$$

$$\mathbf{y}_{n+1} = \mathbf{y}_n + h\mathbf{f}(\mathbf{y}_{n+1})$$

*semi-implicit* linearized eq.

$$\mathbf{y}_{n+1} = \mathbf{y}_n + h \left[ \mathbf{f}(\mathbf{y}_n) + \frac{\partial \mathbf{f}}{\partial \mathbf{y}} \Big|_{\mathbf{y}_n} \cdot (\mathbf{y}_{n+1} - \mathbf{y}_n) \right]$$

only first order accurate!

$$\mathbf{y}_{n+1} = \mathbf{y}_n + h \left[ \mathbf{1} - h \frac{\partial \mathbf{f}}{\partial \mathbf{y}} \right]^{-1} \cdot \mathbf{f}(\mathbf{y}_n)$$

# Higher order implicit methods

## Rosenbrock Methods

$$y(x_0 + h) = y_0 + \sum_{i=1}^s c_i \mathbf{k}_i$$

autonomous differential equation

$$y' = f(y).$$

$$\begin{pmatrix} \mathbf{y} \\ x \end{pmatrix}' = \begin{pmatrix} \mathbf{f} \\ 1 \end{pmatrix}$$

linearized semi-implicit eq.; generalization of embedded RK

$$(\mathbf{1} - \gamma h \mathbf{f}') \cdot \mathbf{k}_i = h \mathbf{f} \left( \mathbf{y}_0 + \sum_{j=1}^{i-1} \alpha_{ij} \mathbf{k}_j \right) + h \mathbf{f}' \cdot \sum_{j=1}^{i-1} \gamma_{ij} \mathbf{k}_j, \quad i = 1, \dots, s$$

we won't go through details, see NR

$$\mathbf{g}_i = \sum_{j=1}^{i-1} \gamma_{ij} \mathbf{k}_j + \gamma \mathbf{k}_i$$

$$(\mathbf{1}/\gamma h - \mathbf{f}') \cdot \mathbf{g}_1 = \mathbf{f}(\mathbf{y}_0)$$

$$(\mathbf{1}/\gamma h - \mathbf{f}') \cdot \mathbf{g}_2 = \mathbf{f}(\mathbf{y}_0 + a_{21} \mathbf{g}_1) + c_{21} \mathbf{g}_1/h$$

$$(\mathbf{1}/\gamma h - \mathbf{f}') \cdot \mathbf{g}_3 = \mathbf{f}(\mathbf{y}_0 + a_{31} \mathbf{g}_1 + a_{32} \mathbf{g}_2) + (c_{31} \mathbf{g}_1 + c_{32} \mathbf{g}_2)/h$$

$$(\mathbf{1}/\gamma h - \mathbf{f}') \cdot \mathbf{g}_4 = \mathbf{f}(\mathbf{y}_0 + a_{41} \mathbf{g}_1 + a_{42} \mathbf{g}_2 + a_{43} \mathbf{g}_3) + (c_{41} \mathbf{g}_1 + c_{42} \mathbf{g}_2 + c_{43} \mathbf{g}_3)/h$$

solve via LU decomp.

# Semi-implicit extrapolation

**Semi-implicit Extrapolation** implicit generalization of Bulirsch-Stoer

implicit mid-pt  $\mathbf{y}_{n+1} - \mathbf{y}_{n-1} = 2h\mathbf{f}\left(\frac{\mathbf{y}_{n+1} + \mathbf{y}_{n-1}}{2}\right)$

$$\Delta_k \equiv \mathbf{y}_{k+1} - \mathbf{y}_k, \quad \Delta_0 = \left[\mathbf{1} - h\frac{\partial\mathbf{f}}{\partial\mathbf{y}}\right]^{-1} \cdot h\mathbf{f}(\mathbf{y}_0)$$

$$\mathbf{y}_1 = \mathbf{y}_0 + \Delta_0$$

even order terms in error => can apply Richardson extrapolation

$k = 1, \dots, m - 1$ , set

$$\Delta_k = \Delta_{k-1} + 2 \left[\mathbf{1} - h\frac{\partial\mathbf{f}}{\partial\mathbf{y}}\right]^{-1} \cdot [h\mathbf{f}(\mathbf{y}_k) - \Delta_{k-1}]$$

$$\mathbf{y}_{k+1} = \mathbf{y}_k + \Delta_k$$

$$\Delta_m = \left[\mathbf{1} - h\frac{\partial\mathbf{f}}{\partial\mathbf{y}}\right]^{-1} \cdot [h\mathbf{f}(\mathbf{y}_m) - \Delta_{m-1}]$$

$$\bar{\mathbf{y}}_m = \mathbf{y}_m + \Delta_m$$

many more methods, predictor-corrector, etc.



# Two-point BVPs

# BVP

$$\frac{dy_i(x)}{dx} = g_i(x, y_1, y_2, \dots, y_N) \quad i = 1, 2, \dots, N$$

$x_1$ , the solution is supposed to satisfy

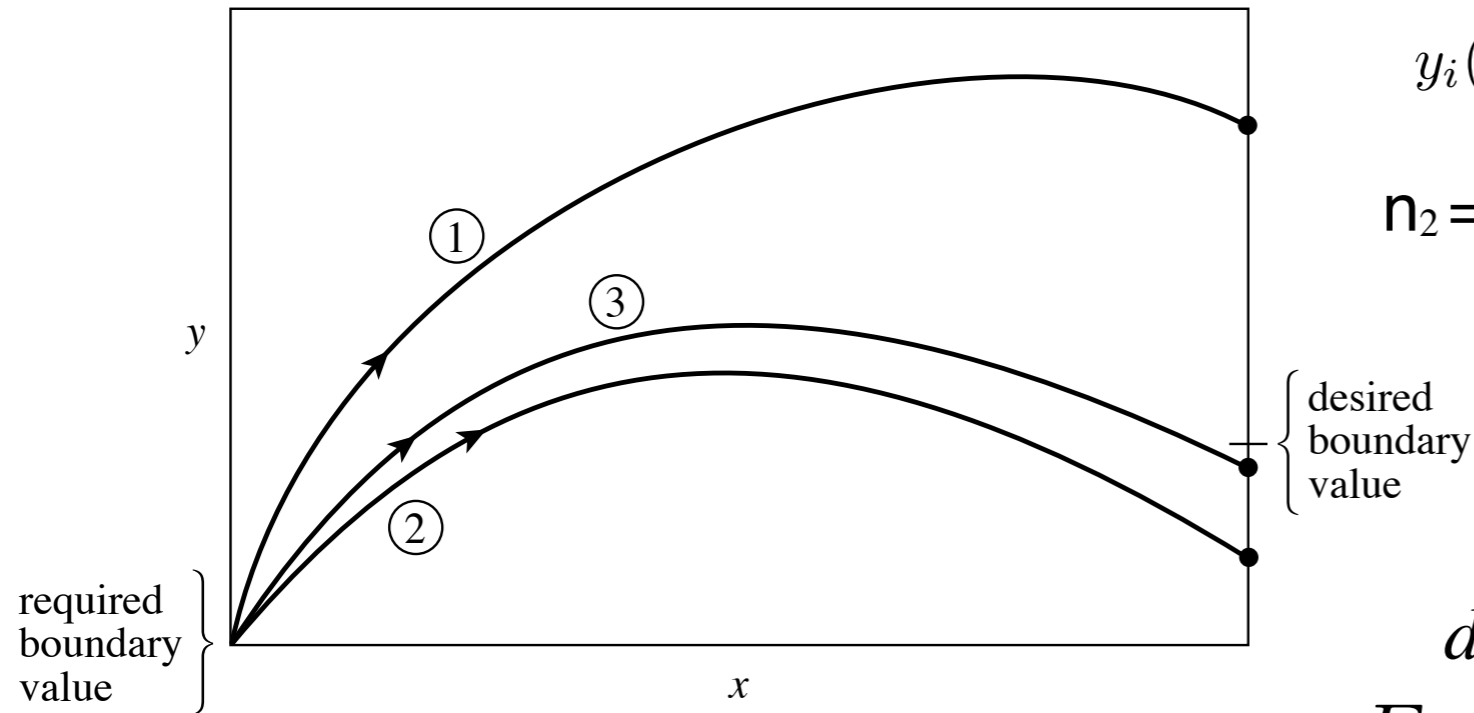
$$B_{1j}(x_1, y_1, y_2, \dots, y_N) = 0 \quad j = 1, \dots, n_1$$

$x_2$ , it is supposed to satisfy

$$n_1 + n_2 = N$$

$$B_{2k}(x_2, y_1, y_2, \dots, y_N) = 0 \quad k = 1, \dots, n_2$$

# Shooting Method



$$y_i(x_1) = y_i(x_1; V_1, \dots, V_{n_2}) \quad i = 1, \dots, N$$

$n_2 = N - n_1$  *freely specifiable* starting values

solve for  $y(x_2)$

*discrepancy vector*  $\mathbf{F}$

$$F_k = B_{2k}(x_2, \mathbf{y}) \quad k = 1, \dots, n_2$$

components of  $\mathbf{F}=0$  only for the desired solution

we want to solve multi-dim. roots of  $\mathbf{F}(V_1, \dots, V_{n_2})=0$ ; use Newton-Raphson

$$\mathbf{J} \cdot \delta \mathbf{V} = -\mathbf{F} \quad \mathbf{V}^{\text{new}} = \mathbf{V}^{\text{old}} + \delta \mathbf{V}$$

$$J_{ij} = \frac{\partial F_i}{\partial V_j} \approx \frac{F_i(V_1, \dots, V_j + \Delta V_j, \dots) - F_i(V_1, \dots, V_j, \dots)}{\Delta V_j}$$

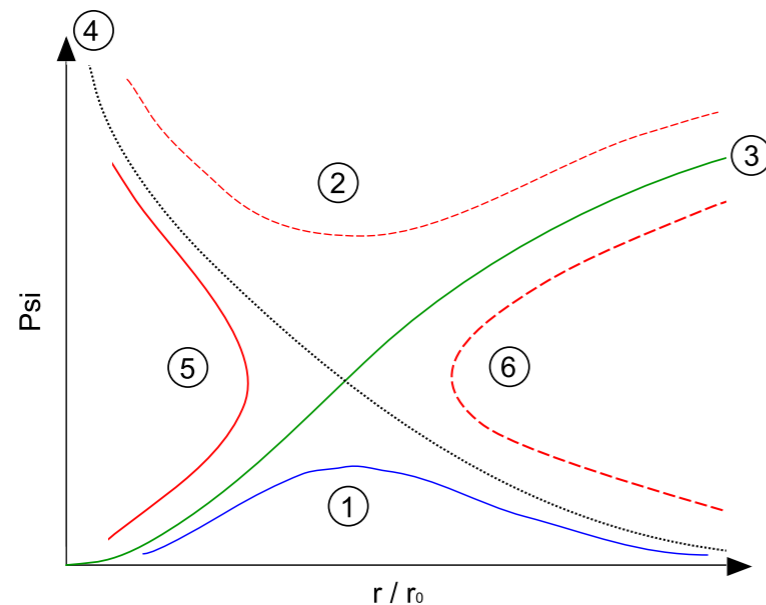
# Shooting to fitting point

useful if singularities, e.g., sonic point

$$y_i(x_1) = y_i(x_1; V_{(1)1}, \dots, V_{(1)n_2}) \quad i = 1, \dots, N$$

$$y_i(x_2) = y_i(x_2; V_{(2)1}, \dots, V_{(2)n_1}) \quad i = 1, \dots, N$$

$$y_i(x_f; \mathbf{V}_{(1)}) = y_i(x_f; \mathbf{V}_{(2)}) \quad i = 1, \dots, N$$



match solution at appropriate fitting point

# Relaxation Methods

$$\frac{dy}{dx} = g(x, y)$$

$$y_k - y_{k-1} - (x_k - x_{k-1}) g \left[ \frac{1}{2}(x_k + x_{k-1}), \frac{1}{2}(y_k + y_{k-1}) \right] = 0 \quad \text{FDE}$$

$$0 = \mathbf{E}_k \equiv \mathbf{y}_k - \mathbf{y}_{k-1} - (x_k - x_{k-1}) \mathbf{g}_k(x_k, x_{k-1}, \mathbf{y}_k, \mathbf{y}_{k-1}), \quad k = 2, 3, \dots, M \quad \text{N(M-1) eqs. at interior pts.}$$

$$0 = \mathbf{E}_1 \equiv \mathbf{B}(x_1, \mathbf{y}_1) \quad n_1 \text{ BCs at } x_1 \quad 0 = \mathbf{E}_{M+1} \equiv \mathbf{C}(x_M, \mathbf{y}_M) \quad n_2 \text{ BCs at } x_M, n_1 + n_2 = N$$

NM nonlinear eqs. for NM unknowns  $y_j^k$ ;  $j=1..N$ ;  $k=1..M$ : can be solved via multi-D Newton-Raphson

Taylor series expansion at interior points;  $k=2..M$ :  $\mathbf{E}_k(\mathbf{y}_k + \Delta\mathbf{y}_k, \mathbf{y}_{k-1} + \Delta\mathbf{y}_{k-1}) \approx \mathbf{E}_k(\mathbf{y}_k, \mathbf{y}_{k-1})$

$$S_{j,n} = \frac{\partial E_{j,k}}{\partial y_{n,k-1}}, \quad S_{j,n+N} = \frac{\partial E_{j,k}}{\partial y_{n,k}}, \quad n = 1, 2, \dots, N$$

$$+ \sum_{n=1}^N \frac{\partial \mathbf{E}_k}{\partial y_{n,k-1}} \Delta y_{n,k-1} + \sum_{n=1}^N \frac{\partial \mathbf{E}_k}{\partial y_{n,k}} \Delta y_{n,k}$$

$$\sum_{n=1}^N S_{j,n} \Delta y_{n,k-1} + \sum_{n=N+1}^{2N} S_{j,n} \Delta y_{n-N,k} = -E_{j,k}, \quad j = 1, 2, \dots, N$$

$$\sum_{n=1}^N S_{j,n} \Delta y_{n,1} = -E_{j,1}, \quad j = n_2 + 1, n_2 + 2, \dots, N$$

inner boundary

$$\sum_{n=1}^N S_{j,n} \Delta y_{n,M} = -E_{j,M+1}, \quad j = 1, 2, \dots, n_2$$

outer boundary

$$S_{j,n} = \frac{\partial E_{j,1}}{\partial y_{n,1}}, \quad n = 1, 2, \dots, N$$

$$S_{j,n} = \frac{\partial E_{j,M+1}}{\partial y_{n,M}}, \quad n = 1, 2, \dots, N$$

start with initial guess  $y_0$ ; add  $\Delta y$  in every iteration till desired accuracy

matrix-eq. for  $n=1..5$  (5 dependent vars.),  $k=1..4$  (4 grid pts.), 3/2 eqs. at inner/outer bdry

X X X X X	V	B
X X X X X	V	B
X X X X X	V	B
X X X X X X X X X X	V	B
X X X X X X X X X X	V	B
X X X X X X X X X X	V	B
X X X X X X X X X X	V	B
X X X X X X X X X X	V	B
X X X X X X X X X X	V	B
X X X X X X X X X X	V	B
X X X X X X X X X X	V	B
X X X X X X X X X X	V	B
X X X X X X X X X X	V	B
X X X X X X X X X X	V	B
X X X X X X X X X X	V	B
X X X X X X X X X X	V	B
X X X X X X X X X X	V	B
X X X X X X X X X X	V	B
X X X X X	V	B
X X X X X	V	B

V: unknown increments  
 B: known  $E_{j,k}$   
 X: known coupling coefficients S  
 block-diagonal matrix  
 Gaussian-Elimination inexpensive  
 see NR for details